



Nonlinear Aeroelastic Optimization

The Cultural and Convention Center
METU

Inonu bulvari
Ankara, Turkey

Sponsored by:
RTA-NATO

The Applied Vehicle Technology Panel

presented by
R.M. Kolonay Ph.D.
General Electric Corporate Research & Development Center
Ankara, Turkey Oct.. 1-5, 2001



Presentation Outline

- Introduction
- Nonlinear Unsteady Aerodynamic Approximations
- Nonlinear Unsteady Aeroelastic Analysis for Design
- Nonlinear Unsteady Aeroelastic Sensitivity Analysis
- Nonlinear Unsteady Aeroelastic Optimization (Transonic)
- Nonlinear Static Aeroelasticity Analysis for Design (Transonic)



Introduction

Goal of Computational Aeroelastic Design Methods

- To accurately *predict* static and dynamic response/stability at a *reasonable computational cost*
- Methods need to be suitable for incorporation into MDA and MDO Environments



Introduction

Discretization of EOM

- Structures K, B, M - Typically, although not necessarily, represented by Finite Elements in either physical or generalized coordinates. Derived in a *Lagrangian* frame of reference.
- External Loads $F(u, \dot{u}, t)$ - Aerodynamic loads. Representations range from Prandtl's lifting line theory to full Navier-Stokes with turbulence modeling. Represented in physical and generalized coordinates in a (usually) *Eulerian* frame of reference.



Introduction

Fluid-Structural Coupling Requirements

- Must ensure spatial compatibility - proper energy exchange across the fluid-structural boundary
- Time marching solutions require proper time synchronization between fluid and structural systems
- For moving CFD meshes GCL [1] must be satisfied

If coupling requirements for time-accurate aeroelastic simulation are not met then dynamical equivalence cannot be achieved. That is, regardless of the fineness of the CFD/CSM meshes and the reduction of time step to 0, the scheme may converge to the “wrong” equilibrium/instability point.[2]



Introduction

Nonlinear Aerodynamics

- Assume a continuum (Conservation of Mass, Momentum, Energy)

Decreasing Degree of Approximation

Complete
No Viscosity
Irrotational, Isentropic
Small Perturbations
Linearize

Equations

Navier-Stokes Equations
Euler Equations
Full Potential Equations
Transonic Small Disturbance
Linear Lifting Surface Theory



Introduction

Nonlinear Flow Conditions

- High angles of attack
- Large control surface deflections
- Transonic Speeds
- . . .

Nonlinear Flow Characteristics

- Attached flows with shocks
- Mixed attached and separated flows
- Mixed attached and separated flows with shocks
- Fully separated flows
- Vortex flows
- . . .



Introduction

General Modeling Comments

- Use appropriate theory to capture desired phenomena
 - Fluids - Navier-Stokes vs. Prandtl's' lifting line theory
 - Structures - Nonlinear FEM vs. Euler beam theory
- Model the fluid and structure with a consistent fidelity
 - For a wing do not model the fluid with NS and the structure with beam theory
- For design methods remember:
 - Total Load = Weight*Nz
 - Stability is determined by a perturbation about a steady state



Introduction

Aeroelastic Phenomena

Static Aeroelastic Phenomena

- Lift Effectiveness
- Divergence
- Control Surface Effectiveness/Reversal
- Aileron Effectiveness/Reversal

Dynamic Aeroelastic Phenomena

- Flutter
- Gust Response
- Buffet
- Limit Cycle Oscillations (LCO)
- Panel Flutter
- Transient Maneuvers
- Control Surface Buzz



Introduction

Equations of Motion

The equations of motion can be expressed in generalized coordinates as

$$[\bar{M}]\{\dot{q}\} + [\bar{B}]\{\dot{q}\} + [\bar{K}]\{q\} = \{\bar{F}\} \quad (1)$$

Where $[\bar{M}]$, $[\bar{K}]$, $[\bar{B}]$, $\{\bar{F}\}$ are generalized mass, stiffness, damping, and aerodynamic forces respectively.

$\{q\}$ - set of independent generalized coordinates defined by $\{u(t)\} = [\Phi]\{q(t)\}$.

$\{u(t)\}$ - spatially discretized structural dof.

$[\Phi]$ - Modal transformation matrix

And the generalized aerodynamic forces can be represented as

$$\bar{F}_i(t) = \frac{\rho_\infty U_\infty^2}{2} C_R^2 \iint_{\hat{S}} [C_{p_L}(x, y, t) - C_{p_U}(x, y, t)] \Phi_i(x, y) dS \quad (2)$$



Unsteady Aerodynamic Approximations

Some Approximation Methods

- Reduced Order Unsteady Aerodynamics
- Pulse Transfer Function (Linear Impulse Function)
- Indicial Response Method
- Discrete-Time Linear and Nonlinear Aerodynamic Impulses (Volterra Theory)



Unsteady Aerodynamic Approximations

Reduced Order Unsteady Aerodynamics [3]

- Use Karhunen-Loeve (K-L) Modes as basis for fluid system

K-L eigenvectors $[V]$ are determined from the eigenvalue problem $[\Phi]\{v\} = \lambda\{v\}$

$$\Phi_{j,k} = \left\{ \tilde{q}_c^j \right\}^T \left\{ \tilde{q}_c^k \right\}$$

$\left\{ \tilde{q}_c^j \right\} \equiv \{ \tilde{q}^j \} - \{ \tilde{q}' \}$ - the deviation of each instantaneous flow field from the mean flow.

$\{ \tilde{q}^j \}$ - an instantaneous flow field vector retained at J discrete times. ($j = 1, 2, 3, \dots, J$)

$\{ \tilde{q}' \}$ - mean flow field.

- Use $[V]$ to reduce the aerodynamic system from $N \times N$ to $R \times R$ where $R \ll N$.
- Demonstrated on 2-D Euler.



Unsteady Aerodynamic Approximations

Pulse Transfer Function [4]

- Determine the response due to a smoothly varying pulse (exponential pulse) for each structural mode participating in the analysis.
- Unsteady generalized aerodynamic forces in *frequency* domain are determined by dividing the Fast Fourier Transform (FFT) of the time domain generalized displacements with the FFT of the time domain generalized aerodynamic forces (Transfer Function)
- Requires one time domain solution for each mode in analysis
- Assumes dynamic response linear about a nonlinear steady state
- Demonstrated on 3-D Euler and Navier-Stokes



Unsteady Aerodynamic Approximations

Discrete-Time Linear and Nonlinear Aerodynamic Impulses (Volterra Theory) [5]

- Volterra theory of nonlinear systems - Any nonlinear system can be modeled as an infinite sum of multidimensional convolution integrals of increasing order.

$$y(t) = h_o + \int_0^{\infty} h_1(t - \tau)u(\tau)d\tau + \int_0^{\infty} \int_0^{\infty} h_2(t - \tau_1)u(\tau_1)u(\tau_2)d\tau_1 d\tau_2 + \dots$$
$$+ \int_0^{\infty} \dots \int_0^{\infty} h_n(t - \tau_1, \dots, t - \tau_n)u(\tau_1) \dots u(\tau_n)d\tau_1 \dots d\tau_n + \dots$$

$y(t)$ - Response of the nonlinear system to $u(t)$

$u(t)$ - An arbitrary input

h_o - Steady state value about which the response is computed

h_1 - First-order kernel (linear unit impulse response)

h_2 - Second-order kernel

h_n - n th order kernel



Unsteady Aerodynamic Approximations

- Kernel Identification
 - h_1 - Impulse response
 - h_n - Responses of the nonlinear system to multiple unit impulses, with the number of impulses applied to the system equal to the order of the kernel of interest
 - h_2 - Response to two unit impulses applied at times t_1 and t_2 .
- Requires more time integrations to determine higher order kernels
- Generalization of Impulse Function Technique
- Approximation *can capture nonlinear effects* that occur during perturbation



Scope of Remaining Discussions

- To develop a transonic unsteady aeroelastic design methodology for *preliminary design* at a reasonable computational cost.
- Such a methodology could be incorporated into the multidisciplinary analysis and design optimization environment (MDAO)



Definition of Preliminary Design

- Geometry assumed fixed (wing planform, airfoil shape, #spars, #ribs, spar spacing etc.)
- Designed components consist of structural elements (thicknesses of skins, ribs and spars, cross sectional areas of posts, spar and rib caps, and concentrated masses etc.)



Critical Issues/Requirements

- Transonic aerodynamics are nonlinear
- Optimization problem is nonlinear
- Coupling two nonlinear problems is not realistic

Preliminary Design MDDA Requirements

- Structural and mass distribution down to substructure level
- F.E.M 50k to 100k d.o.f maximum
- Aerodynamic loads normal to lifting surface
- Small disturbance theory (location and strength of weak shocks)



Solution Approach

- Formulate as nonlinear mathematical programming problem
- Solve by gradient based optimization
- Analysis
 - F.E.M
 - TSD with Indicial Response Method to approximate the unsteady aerodynamic forces
 - * Modal based p -Method for flutter analysis
- Constraint/Constraint gradients
 - Constraint on modal damping rather than on flutter velocity
 - Semi-analytic gradients (some assumptions can produce fully analytic gradients)
- Redesign
 - Use First Order Taylor Series Approximations for functional values and constraints
 - Solve approximate problem by Method of Modified Feasible Directions



Unsteady Aerodynamic Approximations

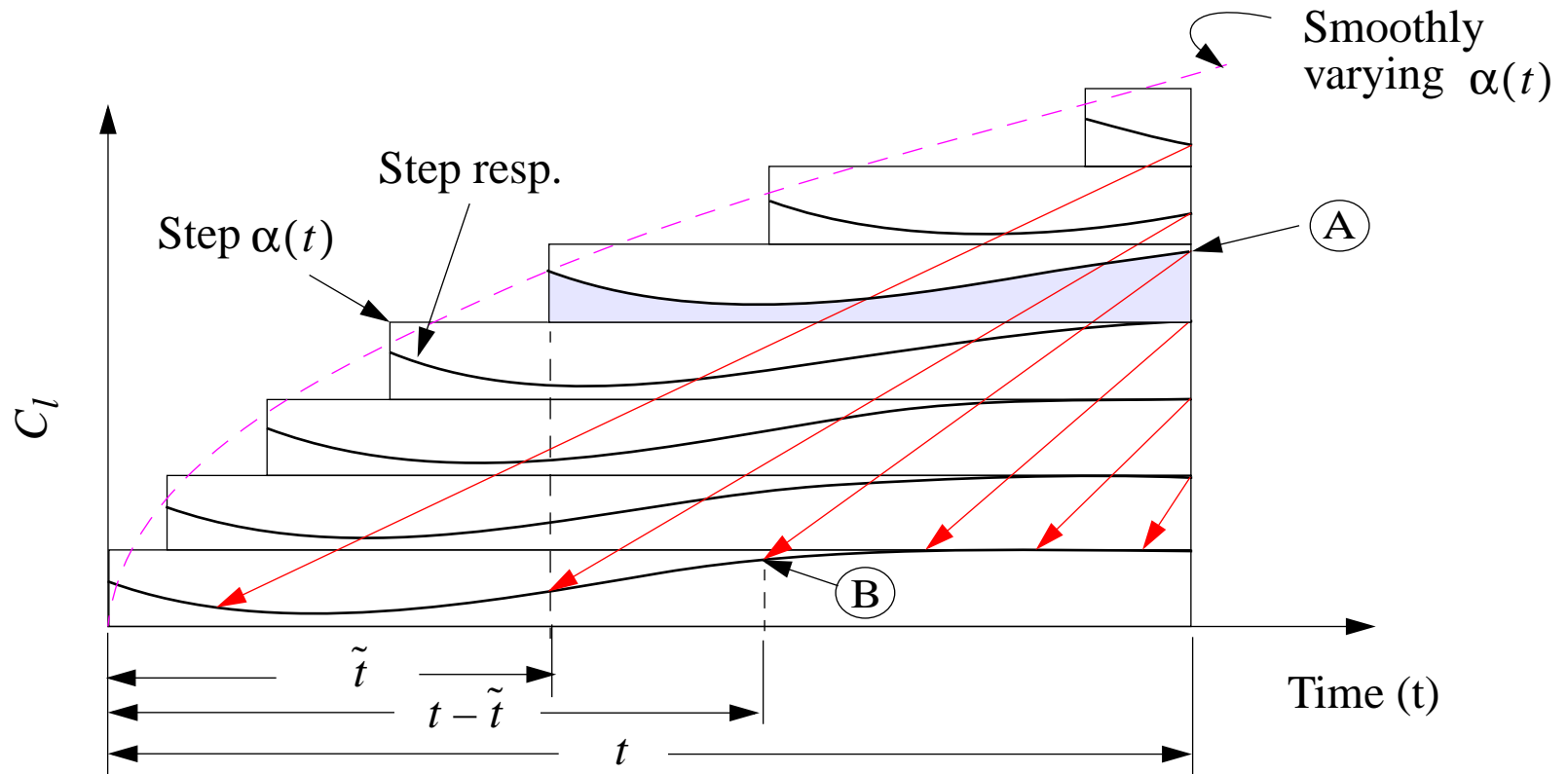
Indicial Response Method [6],[10]

- Assume the unsteady aerodynamic forces are linear about some non-linear static aeroelastic solution (valid for small perturbations).
- Allows use of an *Indicial Response Method* (IRM) to approximate unsteady aerodynamic forces (limited linear superposition).
- IRM requires one nonlinear solution for each mode that participates in flutter analysis for a given set of initial conditions (Mach number, initial angle of attack, Reynolds #, etc.).
- Small perturbations - can neither create, destroy nor significantly move a shock.
- Impulse Function and IRM related (Impulse response derivative of Step response)



Unsteady Aerodynamic Approximations

Indicial Response Method (IRM)



$$C_l(t) = \sum_{i=0}^n C_{l_{\alpha \tilde{t}_i}}(t) \Delta \alpha(\tilde{t}_i) \text{ or } C_l(t) = \sum_0^t C_{l_{\alpha t_0}}(t - \tilde{t}) \Delta \alpha(\tilde{t}) \quad (3)$$



Unsteady Aerodynamic Approximations (IRM)

In the limit $\Delta t_i \rightarrow 0$ and also with a change of variable $\tau = t - \tilde{t}$ Eq (3) becomes

$$C_l(t) = C_{l_{\alpha t_0}}(t)\alpha(0) + \int_0^t C_{l_{\alpha t_0}}(\tau)\dot{\alpha}(t-\tau)d\tau \quad (4)$$

Eq (4) is a form of Duhamel's formula. It enables the calculation of the response of a system to an arbitrary input by knowing only the indicial response of the system. It assumes that linear superposition applies.

Now for a generalized force $\bar{F}_i(t)$, step $q_j(t)$ $Q_{ij}(t)$ Eq can be written as

$$\frac{\bar{F}_i(t)}{\frac{1}{2}\rho_\infty U_\infty^2 C_R^2} = Q_{ij}(t)q_j(0) + \int_0^t Q_{ij}(\tau)\dot{q}_j(t-\tau)d\tau \quad (5)$$

Transforming from the time domain to the Laplace domain via the convolution integral gives

$$\bar{F}_i(s) = \frac{1}{2}\rho_\infty U_\infty^2 C_R^2 s Q_{ij}(s)q_j(s) \quad (6)$$

Eq (6) represents generalized unsteady aerodynamics forces in the *s-domain*



Unsteady Aerodynamic Approximations (IRM)

Defining $[\bar{Q}(s)] = C_R^2 s [Q(s)]$ and transforming the rest of Eq. (1) into the s -domain yields

$$\left[[\bar{M}]s^2 + [\bar{K}] + [\bar{B}]s - \frac{1}{2}\rho_\infty U_\infty^2 [\bar{Q}(s)] \right] \{q(s)\} = 0 \quad (7)$$

The Laplace variable s is in general complex and is defined as $s = \sigma + i\omega$

σ - exponential damping

ω - oscillatory frequency of the system

- Solutions to Eq (7) for flutter can be found by p - k , p , root-locus etc. Provided $\bar{Q}(s)$ is a continuous function of s but need not be harmonic.



Unsteady Aerodynamic Approximations (IRM)

- Assume

$$\bar{Q}_{ij}(s) = sC_R^2 \sum_{r=1}^n \frac{C_{r ij}}{s + b_{r ij}} \quad (8)$$

- Determine the coefficients $C_{r ij}, b_{r ij}$ in the following manner

- Determine the static aeroelastic response for a given flight condition
- Determine $\Delta \bar{F}_{ij}(t)_N$ in the generalized force $\bar{F}_i(t)$ due to a Laplace transformable step input $q_{S_j}(t)$ for the j th mode.
- Once $\Delta \bar{F}_{ij}(t)_N$ is obtained $C_{r ij}, b_{r ij}$ are selected to produce the “best” approxi-

mating curve fit between $\Delta \bar{F}_{ij}(t)_N$ and $\Delta \bar{F}_{ij}(t)_A$ where $\Delta \bar{F}_{ij}(t)_A = \mathfrak{S}^{-1} \left\{ \bar{Q}_{ij}(s) q_{S_j}(s) \right\}$

- with $q_j(s) = A_j \Phi_j \left\{ \int_0^a \left(10 \left(\frac{t}{a} \right)^3 e^{-st} - 15 \left(\frac{t}{a} \right)^4 e^{-st} + 6 \left(\frac{t}{a} \right)^5 e^{-st} \right) dt + \int_a^\infty e^{-st} \right\}$



Unsteady Aerodynamic Approximations (IRM)

Let $s = \omega(\gamma + i)$, define $p \equiv k(\gamma + i)$, $k = (\omega C_R)/(2U_\infty)$, $b = \frac{C_R}{2}$ Eq. (7)

becomes

$$\left[\left(\frac{U_\infty}{b} \right)^2 p^2 [\bar{M}] + \left(\frac{U_\infty}{b} \right) p [\bar{B}] + [\bar{K}] - \frac{1}{2} \rho_\infty U_\infty^2 [\bar{Q}(p)] \right] \{q(p)\} = 0 \quad (9)$$

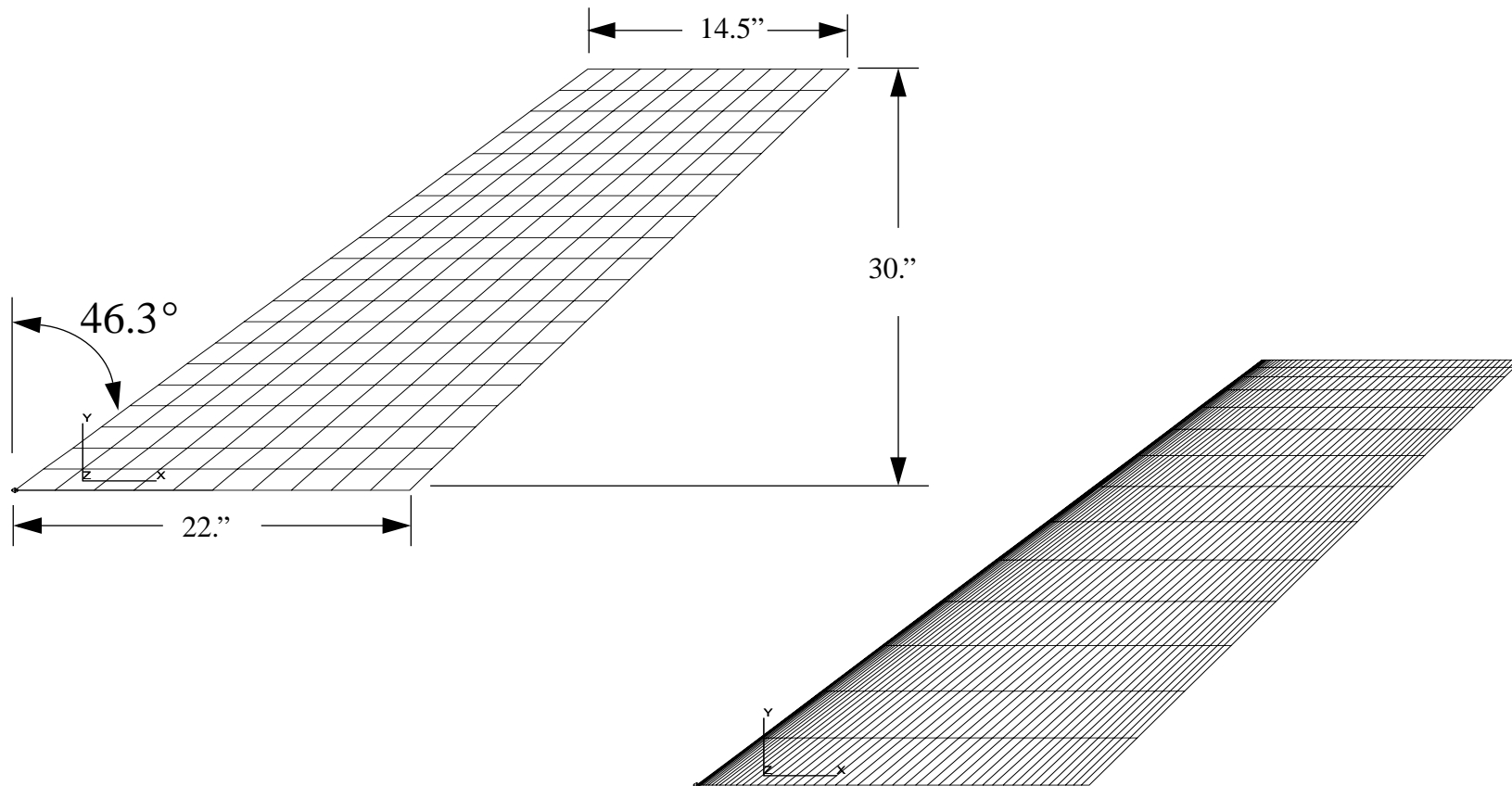
- Solve Eq. (9) by the *p-method*
- *p-method* γ damping valid away from axis (nice for flutter constraints)



Nonlinear Unsteady Aeroelastic Analysis

Indicial Response Flutter Analysis

AGARD 445.6 Wing (Weak Model 3) - TSD Aerodynamics (CAP-TSD)

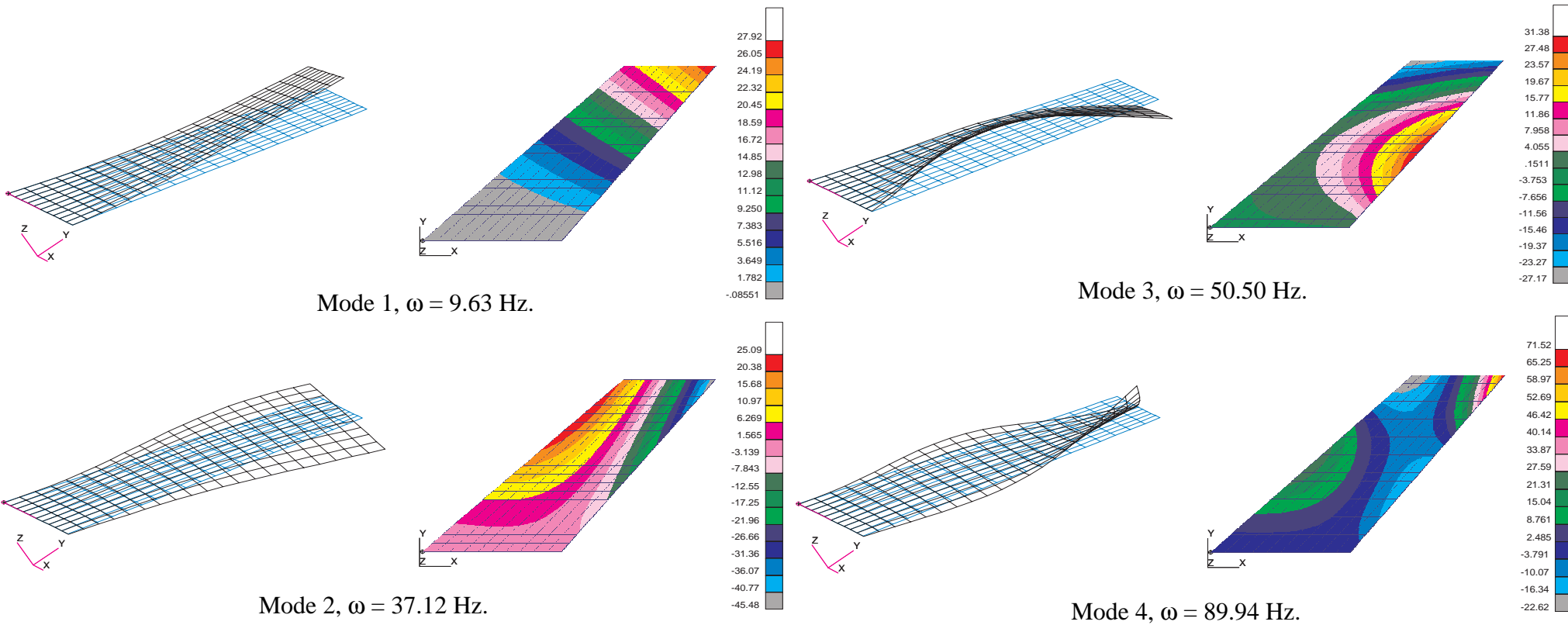




Nonlinear Unsteady Aeroelastic Analysis

445.6 Flutter Analysis

Mode Shapes and frequencies

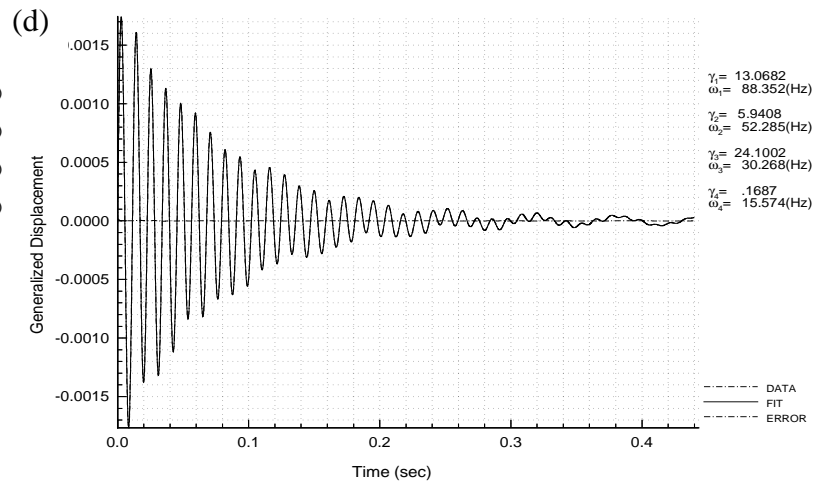
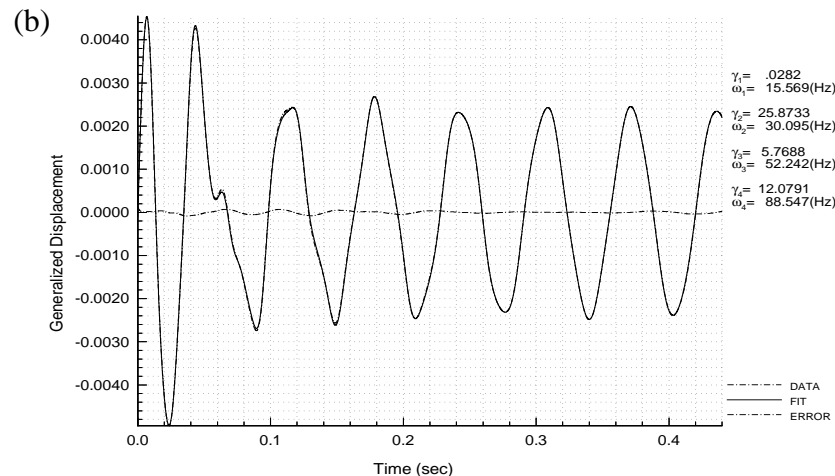
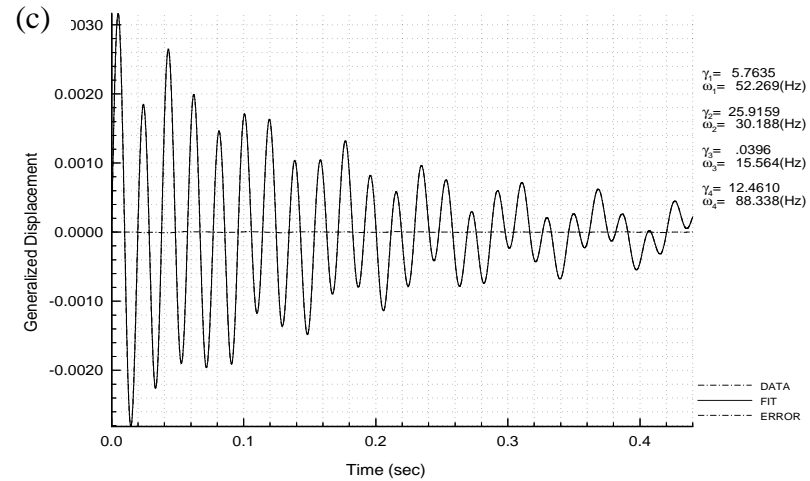
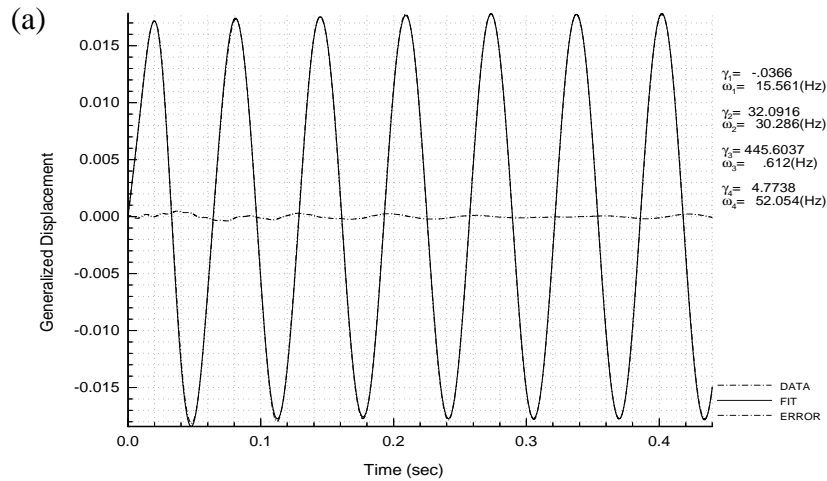




Nonlinear Unsteady Aeroelastic Analysis

445.6 Time Integration Response

$q_f \approx 0.667$ psi, $U_f \approx 11,971$ in/sec, $\omega_f \approx 15.45$ Hz



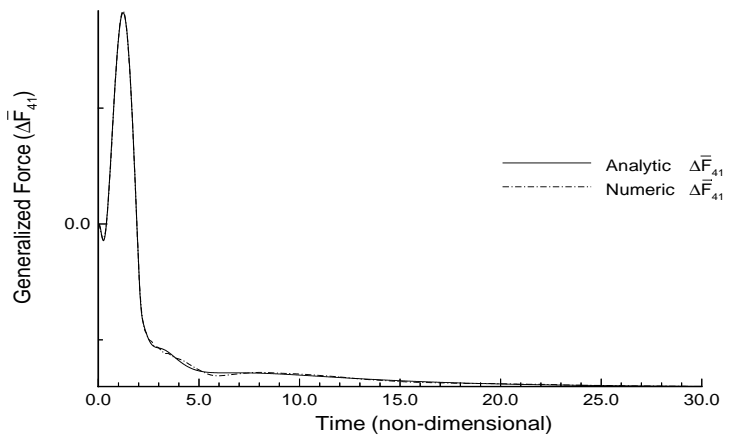
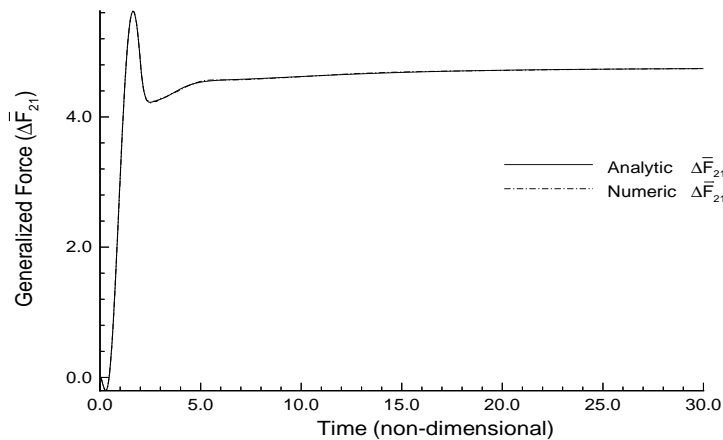
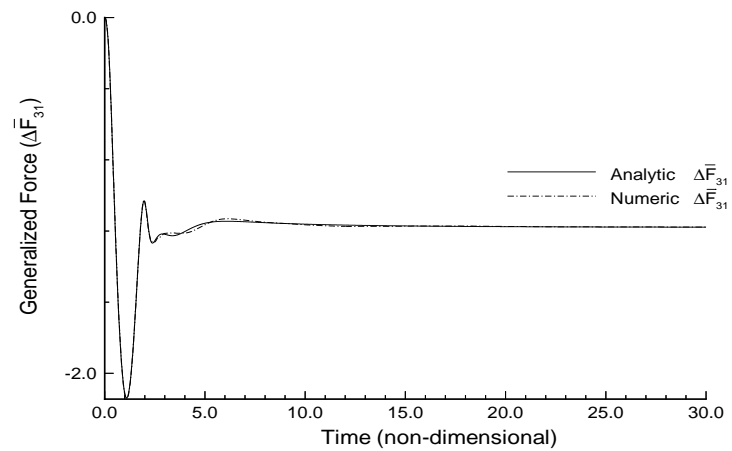
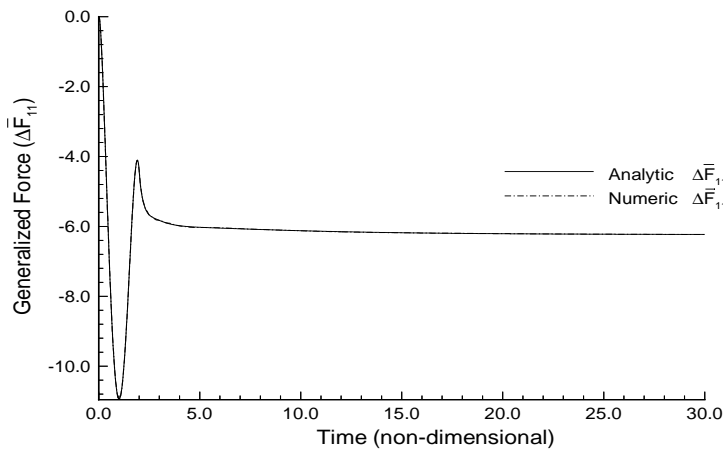
Time response for 445.6 wing (a) mode 1 (b) mode 2 (c) mode 3 (d) mode 4 ($M_\infty = 0.901$, $\alpha_0 = 0.0^\circ$, $q_\infty = 0.67$ psi, $U_\infty = 11998.75$ in/sec)



Nonlinear Unsteady Aeroelastic Analysis

445.6 Indicicial Responses $\Delta\bar{F}_{j1}$

$$M_\infty = 0.901, \alpha_0 = 0.0^\circ$$



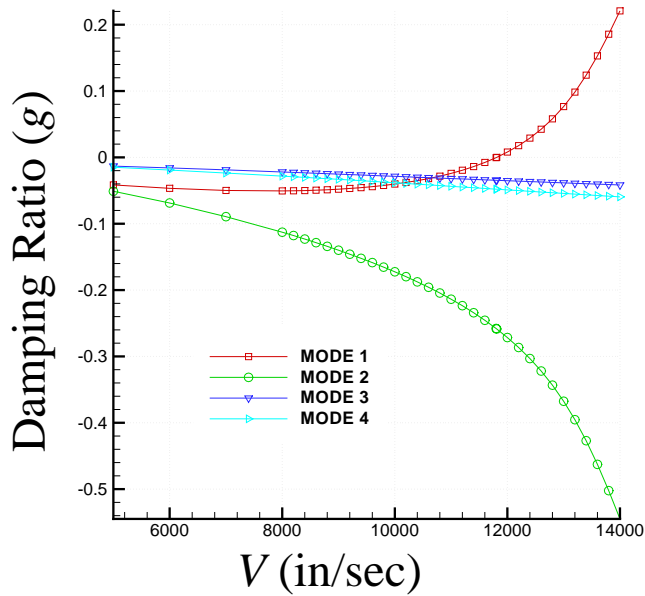


Nonlinear Unsteady Aeroelastic Analysis

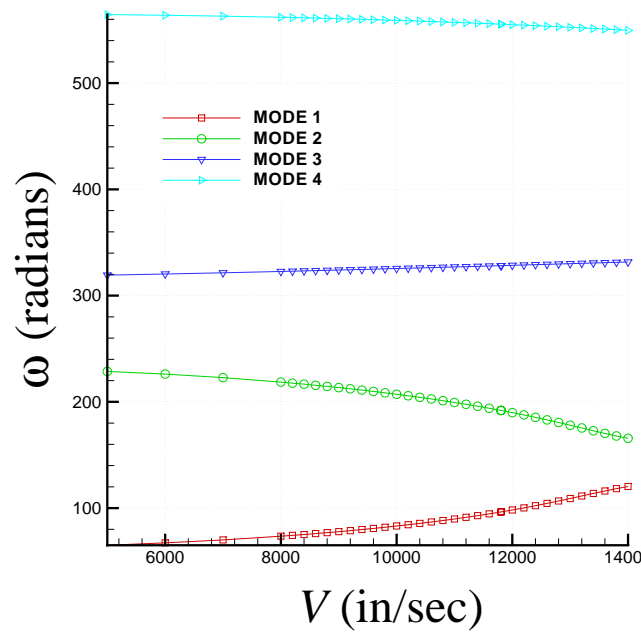
445.6 Indicical Response Flutter Analysis

$$q_f \approx 0.650 \text{ psi}, U_f \approx 11,810 \text{ in/sec}, \omega_f \approx 15.35 \text{ Hz}$$

V vs. g



V vs. ω



g vs. ω

